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Abstract: Nonempirical LCAO-MO-SCF calculations employing a double  $\zeta$  basis set have been carried out to study two cross-sections of the potential surface of methylcarbene. It was found that methylcarbene has a singlet ground state with an  $(S_0-T_1)$  separation of 0.3 kcal/mol. Calculations were also performed for ethylene and various possible transition-state structures through which the rearrangement of methylcarbene to ethylene may occur. The results suggest that this rearrangement involves the migration of a methyl hydrogen gauche to the methine proton.

D ivalent reactive intermediates such as carbenes are of great theoretical interest, <sup>1-4</sup> since it is conceivable that these may exist either in a singlet or in a triplet state.

The simplest carbene, methylene, has been studied extensively.<sup>1,2,5,6</sup> Calculations on this molecule have been reported by several workers,<sup>1,2</sup> all of them predicting a triplet ground state with a singlet-triplet (S<sub>0</sub>-T<sub>1</sub>) separation in the range of 25–37 kcal/mol. Experimental observations<sup>6</sup> are in good agreement with the above. The reactions of methylene have also been studied. Theoretical calculations have been carried out<sup>2,3</sup> for the mechanism and stereochemistry of insertion and addition.

Other carbenes, however, prefer to rearrange via a 1,2-shift of a vicinal hydrogen rather than undergo the above reactions.<sup>7</sup> An example is methylcarbene (1) which rearranges to ethylene (2). This reaction could

$$\begin{array}{c} CH_{3}\ddot{C}H \longrightarrow CH_{2}CH_{2} \\ 1 & 2 \end{array}$$

occur in a variety of ways depending upon the electronic state of the reactant and the state in which the ethylene is produced.<sup>4</sup> The stereochemistry of the above process was studied by Tee,<sup>4</sup> who, on the basis of least motion calculations,<sup>8</sup> predicted that the reaction will most probably involve a triplet methylcarbene rearranging to a "twisted" triplet ethylene. However, least motion calculations consider solely the geometries of the various

(1) For recent *ab initio* calculations, see (a) J. F. Harrison and L. C. Allen, *J. Amer. Chem. Soc.*, **91**, 807 (1969); (b) C. F. Bender and J. F. Shaeffer III, *ibid.*, **92**, 4984 (1970); (c) S. Y. Chu, A. K. O. Siu, and E. F. Hayes, *ibid.*, **94**, 2969 (1972), and references therein.

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(3) R. C. Dobson, D. M. Hayes, and R. Hoffmann, J. Amer. Chem. Soc., 93, 6188 (1971).

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Chem. Soc., 92, 7491 (1970); (c) R. A. Bernheim, H. W. Bernard, P. S.
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(8) For the development of the calculations applying the Principle of Least Motion (PLM), see (a) J. Hine, J. Amer. Chem. Soc., 88, 5525 (1966); (b) O. S. Tee, *ibid.*, 91, 7144 (1969).

reactant and product conformations, and neglect the energy differences between them. Furthermore, prior to this publication, neither the nature and the energy of the ground state of methylcarbene nor the magnitude of separation between its various electronic states had been firmly established. Consequently, the predictions of the PLM approach<sup>8</sup> are valid only if the energy differences between singlet and triplet methylcarbene are smaller than the differences in energies of the various transition states leading to the formation of the products.

Recently, Dewar reported<sup> $\theta$ </sup> the results of MINDO/2 type of calculations on methylcarbene. He also investigated the stereochemistry of hydrogen migration, but considered only one of the possibilities, namely the rearrangement of singlet methylcarbene to ground-state ethylene.

The purpose of our investigation was to study some of the cross-sections of the potential energy hypersurface of methylcarbene and thus examine the energetics of the various stereochemical possibilities of its rearrangement.

### **Computational Details**

Nonempirical LCAO-MO-SCF calculations have been carried out on an IBM 370/165 computer using the IBMOL-IV system.<sup>10</sup> With the exception of the preliminary work, all calculations employed a double  $\zeta$ basis set optimized by Dunning<sup>11</sup> as shown in Table I.

 Table I.
 Contracted<sup>a</sup> Gaussian Basis Set for Hydrogen and Carbon

-Hydrog Ex- ponents	gen s set— Coef- ficients	Carbon Ex- ponents	n s set—— Coef- ficients	Carbo Ex- ponents	on p set Coef- ficients
19.2406 2.8992 0.6534 0.1776	0.032828 0.231208 0.817238 1.000000	4232.6100 634.8820 146.0970 42.4974 14.1892 1.9666 5.1477 0.4962 0.1533	0.002029 0.015535 0.075411 0.257121 0.596555 0.242517 1.000000 1.000000	18.1557 3.9864 1.1429 0.3594 0.1146	0.018534 0.115442 0.386206 0.640089 1.000000

<sup>a</sup> Contracted functions are indicated in boldface type.

<sup>(9)</sup> N. Bodor and M. J. S. Dewar, J. Amer. Chem. Soc., 94, 9103 (1972).

<sup>(10)</sup> A. Veillard, "IBMOL: Computation of Wave Functions for Molecules of General Geometry, Version 4," IBM Research Laboratory, San Jose, Calif.

<sup>(11)</sup> T. H. Dunning, Jr., J. Chem. Phys., 53, 2823 (1970).

 Table II.
 Molecular Parameters Used for

 Calculation of Geometries

	Length,				
Molecule	Bond	Å	Bond an	igle, deg	
Methylcarbene					
H₄					
$H_6 - C_3 - C_2$					
$H_{5}$ $H_{1}$					
Singlet	$C_2 - C_3$	1.48° 1.10°	C.H.H.	96.4	
	$C_3 - H_{4, 5, 6}$	1.104°	$C_2C_3H_{4.4}$	$_{5,6}$ te <sup>b</sup>	
Triplet	$C_2 - C_3$ $C_3 - H_1$	1.46° 1.078ª	C.C.H.	<b>9</b> 0ª	
	C <sub>3</sub> -H <sub>4,5,6</sub>	1.104°	$C_2C_3H_{4,5}$	i, 6 te <sup>b</sup>	
Ethylene					
Ground state					
C=C	C-C	1.337°			
	C-H	1.0861	HCH	117.36	
Excited state					
н					
C—C	C–C	1.63°	HCH	120	
н	C-H	$1.086^{f}$			
Transition state					
Structure I Structure II		For stamic apardinates			
Structure III		see Table III			

<sup>a</sup> This angle was used as a reaction coordinate. Lowest value used is shown above. <sup>b</sup> te = tetrahedral  $\approx 109.5^{\circ}$ . <sup>c</sup> Assumed value based on J. C. D. Brand and D. G. Williamson, *Advan. Phys. Org. Chem.*, **1**, 365 (1963), and ref 9. <sup>d</sup> G. Herzberg and J. W. C. Johns, *J. Chem. Phys.*, **54**, 2276 (1971). <sup>e</sup> W. Haugen and M. Traetteberg, *Acta Chem. Scand.*, **20**, 1726 (1966). <sup>f</sup> H. C. Allen, Jr., and E. L. Plyler, *J. Amer. Chem. Soc.*, **80**, 2673 (1958).

For the preliminary calculations on methylcarbene the SCF MO were expanded in terms of a minimal basis set which in turn was contracted to form a set of primitive Gaussian type functions suggested by Klessinger.<sup>12</sup>

The energies of the excited states were calculated by the virtual orbital technique<sup>13</sup> using

$${}^{1}E_{1} = E_{0} + (\epsilon_{a} - \epsilon_{b}) - J_{ba} + 2K_{ba}$$
$${}^{3}E_{1} = E_{0} + (\epsilon_{a} - \epsilon_{b}) - J_{ba}$$

where  $J_{ba}$  and  $K_{ba}$  are the two-electron integrals over the molecular basis and a and b are the virtual and occupied MO involved in the excitation, respectively.

### **Results and Discussion**

Conformation of Methylcarbene. The two most important internal coordinates of methylcarbene are



 $\alpha$ , the C<sub>3</sub>C<sub>2</sub>H<sub>1</sub> bond angle leading to the in-plane inversion about the methine carbon, and the dihedral angle  $\theta$ , responsible for the torsional mode of motion along the C-C bond. Consequently, our theoretical study involved the examination of the two cross-sections  $E = E(\alpha)$  and  $E = E(\theta)$  of the rotation-inversion potential surface  $E = E(\alpha, \theta)$  of methylcarbene.



Figure 1. Variation of total energy as a function of the CCH bond angle,  $\alpha$  (in-plane inversion potential), for the ground (S<sub>0</sub>), first four excited singlet (S<sub>i</sub>), and triplet (T<sub>i</sub>) states of methyl-carbene as computed using a minimal basis set.

For the first cross-section the internal coordinate was chosen to be the bond angle  $\alpha$ , and a series of preliminary calculations were carried out to generate  $E = E(\alpha)$  using a minimum size basis set. The geometrical parameters used for the calculations are listed in Table II. The  $E(\alpha)$  plot for the closed shell ground  $(S_0)$  state and the first four excited singlet  $(S_i)$ and triplet  $(T_i)$  states is shown in Figure 1. It is apparent that methylcarbene, like methylene itself, prefers a bent geometry. In the  $S_0$  state the minimum on the curve corresponds to a conformation having an  $\alpha$  value of  $107.8^{\circ}$ , whereas in the T<sub>1</sub> state the optimal bending angle is  $130^{\circ}$ . The curve associated with the S<sub>1</sub> state appears to be quite flat with a shallow minimum at an  $\alpha$  value of 140°. It is interesting to note that the curves for both the S<sub>2</sub> and T<sub>2</sub> states have their minima at  $\alpha =$ 180°, meaning that the molecule in these states favors a linear geometry. According to Figure 1 methylcarbene should have a triplet ground state with an adiabatic  $(S_0-T_1)$  separation of 21.6 kcal/mol. The first excited singlet state  $(S_1)$  also lies fairly low, only 24.9 kcal/mol above the  $S_0$  state.

For further refinement of the preliminary work, the calculations were repeated using a double  $\zeta$  atomic basis set. Our results are summarized in Table III and illustrated in Figure 2. As was expected, the general appearance of the curves is very similar to that obtained from the minimal basis set (Figure 1), although the  $\alpha$  values of the lowest energy conformations in the various states are slightly shifted toward larger bond angles.<sup>14</sup>

<sup>(12)</sup> M. Klessinger, Theor. Chim. Acta, 15, 353 (1969).

<sup>(13)</sup> C. C. J. Roothaan, Rev. Mod. Phys., 23, 161 (1951).

<sup>(14)</sup> The minimum on the curve obtained for the S<sub>0</sub> state appears at  $\alpha = 111.8^{\circ}$ , on the curve for the T<sub>1</sub> state at  $\alpha = 133.5^{\circ}$ , and on the curve for the S<sub>1</sub> state at  $\alpha = 144.4^{\circ}$ .

	States				
$\beta$ , deg	<b>S</b> <sub>0</sub>	$S_1/T_1$	$S_2/T_2$	$S_3/T_3$	$S_4/T_4$
100	-77.8881039	-77.800210	-77.493197	-77.518667	-77.459202
		-77.862631	-77.512250		- 77.479219
110	-77.8928781	-77.819898	-77.519182	-77.518808	
			77.534063	-77.537123	-77.501902
120		-77.832935		-77.518654	
		77.889537	-77,552598	77.531600	-77.514703
130	-77.8883572	77.840220	-77.553532	-77.517334	77.483674
		-77.893438	-77.569013	-77.527185	- 77.516839
140	-77.8729344	-77.843143		-77.514887	-77.478358
				-77.523657	-77.511982
150	-77.8615764	-77.842962		-77.511077	
		-77.889021		-77,519773	
179	-77.8424125	-77.838201	-77.580515	-77.502030	
			-77.612561	77.511521	77,467810
215.6	-77.8688922	-77.842621		-77.512716	
			— 77 . 587796	-77.521885	-77.487865
226.5	-77.8814421	77.840953	-77.554797	-77.516244	-77.462622
		77,895926	-77.572105	-77.525880	-77,485905
248.2	-77.8964109				

<sup>a</sup> Where  $\beta = (360 - \alpha)$ . For definition of  $\alpha$  see Figure 2. <sup>b</sup> Calculations employed a double  $\zeta$  basis set.



Figure 2. Variation of total energy as a function of the CCH bond angle,  $\alpha$  (in-plane inversion potential), for the ground (S<sub>0</sub>), first four excited singlet (S<sub>i</sub>), and triplet (T<sub>i</sub>) states of methylcarbene as computed using a double  $\zeta$  basis set.

However, the relative energies of the two lowest states are significantly different from the preliminary results. The adiabatic  $(S_0-T_1)$  separation decreased to 0.3 kcal/ mol, with the  $S_0$  state being the lowest energy state when the methylcarbene is in its most stable staggered conformation. On the other hand, when the molecule assumes its less stable eclipsed conformation, the order of the stabilities of the  $S_0$  and  $T_1$  states reverses to a similar extent.<sup>15</sup> A crossover of the two states occurs at an  $\alpha$  value of 121°.

These results suggest that methylcarbene may have either a singlet or a triplet ground state depending on



Figure 3. Variation of total energy as a function of the C–C torsional angle,  $\theta$  (rotation-potential), for the ground (S<sub>0</sub>), first excited singlet (S<sub>1</sub>), and triplet (T<sub>1</sub>) states of methylcarbene.

whether it assumes a staggered conformation or it is forced into an eclipsed conformation by means of some geometrical constraints. Furthermore, the close energy gap between the two states suggests that the generation of any one of the pure spin states could be an extremely demanding experimental task.<sup>16</sup>

For the other cross-section the internal coordinate was chosen to be the dihedral angle  $\theta$ .<sup>17</sup> The results obtained in the form of  $E = E(\theta)$  for the three lowest states are shown in Figure 3 and the actual values are

(17) The dihedral angle *ijkl* is defined as the angular displacement of kl relative to *ij*, measured counterclockwise along the direction  $k \rightarrow j$ .

<sup>(15)</sup> The adiabatic  $(S_0\text{--}T_1)$  gap is 0.45 kcal/mol for the eclipsed conformation.

<sup>(16)</sup> As a referee has pointed out, it is probable that correlation energies in singlet states are somewhat higher than in triplet states, which would result in an overestimation of the  $(S_0-T_1)$  separation by the SCF method. Furthermore, a recent very refined theoretical study of methylene (including differential correlation energy) revealed that addition of d orbitals to the basis set decreases the  $(S_0-T_1)$  gap substantially (C. F. Bender, H. F. Schaefer III, D. R. Franceschetti, and L. C. Allen, J. Amer. Chem. Soc., 94, 6888 (1972)). Thus an alternative interpretation of our results is that the methylcarbene ground state is most probable ably a singlet.

**Table IV.** Variation of Total Energy with the  $H_6C_3C_2H_1$ Dihedral Angle,  $^{a} \theta$ , for the Ground (S<sub>0</sub>), First Excited Singlet  $(S_1)$ , and Triplet  $(T_1)$  States of Methylcarbene

State	Dihedral angle $\theta$ , deg	Total energy, hartree		
S <sub>0</sub>	0			
	30	77 . 8946793		
	55	77.8963530		
	60	77.8964109		
$T_1$	0			
	30	-77.895041		
	55	-77.895861		
	60	-77.895926		
S <sub>1</sub>	0	-77.843446		
-	30	77.843061		
	55	-77.842711		
	60	-77 842621		

<sup>a</sup> For the definition of the dihedral angle see text.

summarized in Table IV. The curves are symmetrical, closely analogous to the well-known threefold rotational potential curves of ethane. The calculated rotational barrier for the S<sub>0</sub> state is 2.14 kcal/mol, for the T<sub>1</sub> state 1.38 kcal/mol, and for the S<sub>1</sub> state -0.52kcal/mol, the latter value indicating that the eclipsed conformation in that state is slightly more stable.

A summary of our results for methylcarbene is shown in Table V.

(carbene) may be calculated. The values obtained are

## $\Delta H_{\rm f}({\rm carbene}) = [(E_{\rm SCF}({\rm carbene}) -$

## $E_{\rm SCF}(\text{ethylene})] + \Delta H_{\rm f}(\text{ethylene})$

included in Table V. The calculated heat of formation for methylcarbene in the S<sub>0</sub> state is 84.8 kcal/mol but an experimental value is not available for comparison. However, the heat of formation of methylene was reported to be in the range of 91.9<sup>21</sup>-94.6<sup>22</sup> kcal/mol. The difference between these and our estimate, i.e.,  $\Delta \Delta H_{\rm f} = \Delta H_{\rm f}({\rm HCH}) - \Delta H_{\rm f}({\rm CH}_{\rm 3}{\rm CH}) = 7.1-9.8$ kcal/mol, agrees remarkably well with the average  $\Delta\Delta H_{\rm f}$ of methyl substitution in olefins (e.g.,  $\Delta H_{\rm f}(\rm CH_2=CH_2)$ )  $\Delta H_{\rm f}(\rm CH_3CH=CH_2) = 7.9 \ \rm kcal/mol^{20}$ ). On the other hand, as Dewar pointed out<sup>23</sup> such comparisons are truly applicable only to molecules with localized bonds. He reported<sup>9</sup> the calculated heat of formation of singlet methylcarbene as being 66.0 kcal/mol. The large discrepancy between his value and that estimated by the localized bond approach was attributed to "stabilization of methylcarbene by hyperconjugation."9 It was shown<sup>24</sup> that double  $\zeta$  basis set calculations usually predict the heats of formation within 10 kcal/mol of the experimental value. In the present case this is undoubtedly due to the unavoidable neglect of  $\Delta E_{\rm corr}$ in the thermochemical equation above. Therefore it is conceivable that our estimate is somewhat higher than the correct value.

Table V. Summary of the Results of Calculations for Methylcarbene and Ethylene

<u> </u>	Methylcarbene			Ethvlene		
	S <sub>0</sub>	Tı	Sı	$\mathbf{S}_0$	Ťı	Sı
$\angle H_1C_2C_3$ , deg	111.8	133.5	144.4	121.35	120	120
Total energy, hartree	-77.8964109	-77.895926	-77.843446	-78.0115556	-77.910064	-77.716596
Rotational barrier, kcal/mol	2,14	1.38	$-0.52^{a}$		63.7°	185.2°
Heat of formation $\Delta H_{\rm f}$ , kcal/mol	84.8	85.1	117.9	12.5%		

<sup>a</sup> The negative value indicates that the eclipsed conformation is more stable than the staggered conformation. <sup>b</sup> Value taken from literature.<sup>19</sup> <sup>c</sup> The values quoted refer to the barrier to cis-trans isomerization.

Molecular Rearrangement of Methylcarbene. The second part of our study involved the stereochemistry of the rearrangement of methylcarbene to ethylene. In order to be able to construct a correlation diagram for this process, the energies of both the ground and excited states of ethylene had to be calculated. The molecular parameters used for the calculations are listed in Table II18 and the energies obtained are reported in Table V. The energy value obtained for ground-state ethylene was then used to calculate the heat of formation of methylcarbene by means of the following equation.<sup>19</sup> Since  $\Delta H_{\rm f}$ (ethylene) is well

 $\Delta H_{\rm f}({\rm carbene}) - \Delta H_{\rm f}({\rm ethylene}) \approx$ 

$$E_{\rm SCF}({\rm carbene}) - E_{\rm SCF}({\rm ethylene})$$

known<sup>20</sup> and the quantities on the right-hand side of the equation are available from our calculations,  $\Delta H_{f}$ -

The correlation diagram, constructed for methylcarbene and ethylene, is shown in Figure 4. The diagram reveals that the first excited state  $(T_1)$  of ethylene lies 8.6 kcal/mol below the ground state  $(S_0)$  of methylcarbene. Considering the finding that the  $S_0$  and  $T_1$ states of methylcarbene are almost degenerate, two alternative mechanisms may be envisioned for its rearrangement: mode 1, singlet (S<sub>0</sub>) methylcarbene rearranging to ground-state  $(S_0)$  ethylene, or mode 2, triplet  $(T_1)$  methylcarbene rearranging to form initially an excited triplet  $(T_1)$  ethylene, which in turn would decay to the ground  $(S_0)$  state product. The former process may involve the migration of a hydrogen being either syn or anti to the carbenic hydrogen. In the case of the triplet-triplet conversion the anti process is expected to be more favorable since the migration to a "twisted" ethylene need not be accompanied by any torsion of the C-C bond. Although the former route is thermodynamically favored, it is not unlikely that the

<sup>(18)</sup> For the ground-state conformation  $D_{2h}$  symmetry, and for the excited states C<sub>s</sub> symmetry, was used.

<sup>(19)</sup> The equivalence below holds only approximately because  $\Delta E_{\rm corr}$ is not added to the right-hand side of the equation. This is equivalent to assuming that the correlation energy of methylcarbene is identical

<sup>(20) &</sup>quot;Selected Values of Physical and Thermodynamic Properties of Hydrocarbons and Related Compounds," American Petroleum Institute, Research Project 44, Carnegie Press, Pittsburgh, Pa., 1953.

<sup>(21)</sup> W. A. Chupka and C. Lifshitz, J. Chem. Phys., 48, 1109 (1968).
(22) W. A. Chupka, J. Chem. Phys., 48, 2337 (1968).
(23) M. J. S. Dewar, "The Molecular Orbital Theory of Organic Chemistry," McGraw-Hill, New York, N. Y., 1969, p 141.

<sup>(24)</sup> A. C. Hopkinson, K. Yates, and I. G. Csizmadia, Theor. Chim. Acta, 23, 369 (1972).



Figure 4. Correlation diagram of methylcarbene and ethylene as computed employing a double  $\zeta$  basis set.





<sup>*a*</sup> For definition see text.

latter mode is faster if its barrier to the conversion  $(T_1-T_1')$  is lower.

The barrier for the triplet-triplet conversion (mode 2) was calculated in the following way: a transition-state geometry was assumed by allowing the atoms of the molecule to move synchronously to a position halfway between their initial and final spatial position. This was achieved by averaging all the bond lengths and bond angles between reactant and product. The position of the migrating hydrogen was then optimized as shown in



Figure 5. Optimization of the position of the migrating hydrogen [defined as the distance, m, of the migrating hydrogen from the center of the C-C bond] in the various possible transition-state structures considered for the rearrangement of methylcarbene to ethylene.

Figure 5, curves Ia and Ib. The conformation having minimum energy (structure I) has a  $CCH_{migr}$  angle of 49.4°. Similar types of calculations were carried out for each of the transition-state structures (II and III) involved in the singlet-to-ground state conversion as illustrated in Figure 5. Curve II corresponds to the S<sub>0</sub> state of the transition state conformation for syn migration of mode 1 and curve III to that for the anti migration of mode 1.<sup>25</sup> The corresponding optimum conformations were found to have CCH<sub>migr</sub> angles of 54.3 and 52.2° for structure II (syn) and III (anti), respectively. The atomic coordinates of the optimized transition state structures (I, II, and III) are given in Table VI.

The diagram summarizing the energetics of the rearrangement is shown in Figure 6. It is apparent that the rearrangement of singlet methylcarbene to groundstate ethylene (mode 1) is preferred over the triplettriplet route (mode 2) both kinetically and thermodynamically. For the former process, of the two stereochemical alternatives, the syn migration is predicted to be more facile than the anti migration because of the low barrier (20.9 kcal/mol) involved. This is somewhat surprising since intuitively it might have been expected that the hydrogen would preferentially migrate toward the lobe with a lone pair rather than toward the initially empty p orbital on carbon. But if one considers that, as the rearrangement proceeds, the electron density of the  $\sigma$  bond (C-H<sub>migr</sub> bond) progressively diffuses to all three atoms (*i.e.*,  $CH_{m igr} C$ ) involved in the migration, then the preference for mode 2 is not at all unlikely.<sup>26</sup>

(27) H. E. Zimmerman, Accounts Chem. Res., 5, 393 (1972).

<sup>(25)</sup> For these two only the  $S_0$  states are shown since they were the lowest in energy in the manifold.

<sup>(26)</sup> The above rationalization is qualitatively in accord with the mixing of the orbitals in Zimmerman's MO Following approach,<sup>21</sup> as applied to this rearrangement.

Our results agree with those of Dewar,<sup>9</sup> who also found the syn migration to be favored over the anti migration. In contrast, the triplet-triplet conversion was found to be slightly preferred by least motion considerations.<sup>7</sup> It seems therefore that the rearrangement is non-leastmotion-controlled.

## Conclusions

Our results were arrived at by a limited optimization procedure regarding the geometries of the various  $C_2H_4$ species under consideration. A complete optimization of these and in particular the transition-state geometries could change the relative energies by several kcal/mol. This is not expected to influence the overall stereochemical consequences of the rearrangement which involves the *relative ordering* of the energies of the individual species. On the other hand, since optimization generally lowers the energy, it is likely that the barriers to the migrations would decrease substantially. Consequently, the barrier to the preferred syn mode should be regarded as an upper limit, bearing in mind that the reaction is known to be extremely facile.

As for methylcarbene, it was shown that increasing the basis set results in the decrease of the singlet-triplet separation from 21.6 kcal/mol to approximately zero. A very recent, and probably one of the most refined, theoretical studies of methylene, by Bender, *et al.*,<sup>16</sup> resulted in similar conclusions. In our case this decrease warrants caution in the assignment of the ground electronic state. It is suggested that methylcarbene



Figure 6. A summary of thermodynamic and kinetic stabilities of selected low-lying states of various  $C_2H_4$  species.

most probably has a singlet ground state, with the reservation that for a better understanding of the  $(S_0-T_1)$  spacing a complete study of the energy surface is necessary.

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# Stereochemistry of the Thermal Addition of $\beta$ -Pinene to Maleic Anhydride

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**Abstract:** Two stereochemical aspects of the thermal ene reaction between  $\beta$ -pinene and maleic anhydride have been investigated, the configuration of the allylic hydrogen transferred in the reaction and the preference for endo *vs.* exo orientation. Two routes are described to a stereospecifically labeled  $\beta$ -pinene in which the C-3 hydrogen trans to the *gem*-dimethyl bridge is replaced by deuterium. Use of this material in ene reactions with maleic anhydride and dimethyl acetylenedicarboxylate showed that  $95 \pm 5\%$  of the deuterium is transferred. Determination of the absolute configuration at the new asymmetric center in the adduct as *R*, by degradation to (S)-(+)-3methylpentanoic acid, revealed that the major adduct is **2b**, formed by endo addition. Of four possible transitionstate orientations A-D, C is established as the preferred pathway.

The thermal addition of an alkene to another olefin possessing an allylic hydrogen, the so-called "ene" reaction, is one of the simplest reactions of organic chemistry.<sup>1</sup> Though radical<sup>2</sup> and other mechanisms<sup>3</sup> have been advanced, the addition is usually considered to proceed in a symmetry-allowed concerted process<sup>4,5</sup> via a six-membered cyclic transition state (eq 1), unless prohibited by steric factors.<sup>6</sup> Consistent with the concerted mechanism are two important stereochemical observations: (a) the new C-C and C-H bonds are formed cis,<sup>7</sup> and (b) asymmetric induction occurs in the adduct

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